

Separation of Variables – Oscillating Strings

Bernd Schröder

Separation of Variables

1. Solution technique for partial differential equations.

Separation of Variables

1. Solution technique for partial differential equations.
2. If the unknown function u depends on variables x, t , we assume there is a solution of the form $u = X(x)T(t)$.

Separation of Variables

1. Solution technique for partial differential equations.
2. If the unknown function u depends on variables x, t , we assume there is a solution of the form $u = X(x)T(t)$.
3. The special form of this solution function allows us to replace the original partial differential equation with several ordinary differential equations.

Separation of Variables

1. Solution technique for partial differential equations.
2. If the unknown function u depends on variables x, t , we assume there is a solution of the form $u = X(x)T(t)$.
3. The special form of this solution function allows us to replace the original partial differential equation with several ordinary differential equations.
4. Key step: If $f(x) = g(t)$, then f and g must be constant.

Separation of Variables

1. Solution technique for partial differential equations.
2. If the unknown function u depends on variables x, t , we assume there is a solution of the form $u = X(x)T(t)$.
3. The special form of this solution function allows us to replace the original partial differential equation with several ordinary differential equations.
4. Key step: If $f(x) = g(t)$, then f and g must be constant.
5. Solutions of the ordinary differential equations we obtain must typically be processed some more to give useful results for the partial differential equations.

Separation of Variables

1. Solution technique for partial differential equations.
2. If the unknown function u depends on variables x, t , we assume there is a solution of the form $u = X(x)T(t)$.
3. The special form of this solution function allows us to replace the original partial differential equation with several ordinary differential equations.
4. Key step: If $f(x) = g(t)$, then f and g must be constant.
5. Solutions of the ordinary differential equations we obtain must typically be processed some more to give useful results for the partial differential equations.
6. For the equation in this presentation, Fourier series will allow us to get the actual solution of the problem.

The Equation and the Initial and Boundary Conditions

The Equation and the Initial and Boundary Conditions

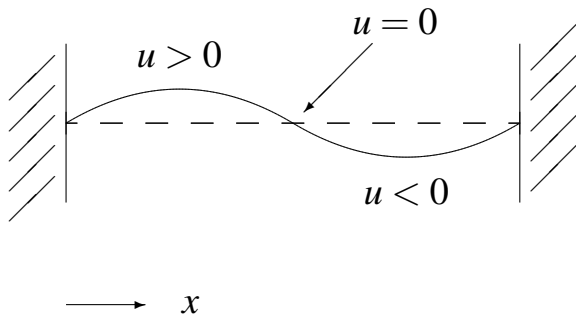
1. $u(x, t)$ is the vertical displacement of the particle at position x on the string,

The Equation and the Initial and Boundary Conditions

1. $u(x, t)$ is the vertical displacement of the particle at position x on the string, at time t .

The Equation and the Initial and Boundary Conditions

1. $u(x, t)$ is the vertical displacement of the particle at position x on the string, at time t .



The Equation and the Initial and Boundary Conditions

1. $u(x, t)$ is the vertical displacement of the particle at position x on the string, at time t .

2. Partial differential equation:
$$\frac{\partial^2}{\partial t^2} u(x, t) = \frac{\partial^2}{\partial x^2} u(x, t).$$

The Equation and the Initial and Boundary Conditions

1. $u(x, t)$ is the vertical displacement of the particle at position x on the string, at time t .
2. Partial differential equation: $\frac{\partial^2}{\partial t^2}u(x, t) = \frac{\partial^2}{\partial x^2}u(x, t)$.
(We're assuming $k = 1$.)

The Equation and the Initial and Boundary Conditions

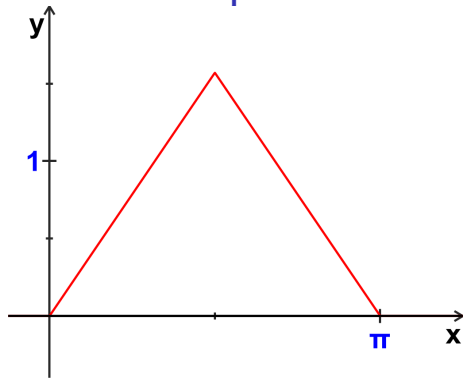
1. $u(x, t)$ is the vertical displacement of the particle at position x on the string, at time t .
2. Partial differential equation: $\frac{\partial^2}{\partial t^2} u(x, t) = \frac{\partial^2}{\partial x^2} u(x, t)$.
(We're assuming $k = 1$.)
3. Initial condition. For $0 \leq x \leq \pi$: $u(x, 0) = f(x)$ (see next slide), $\frac{\partial}{\partial t} u(x, 0) = 0$ (string is not moving as it is released).

The Equation and the Initial and Boundary Conditions

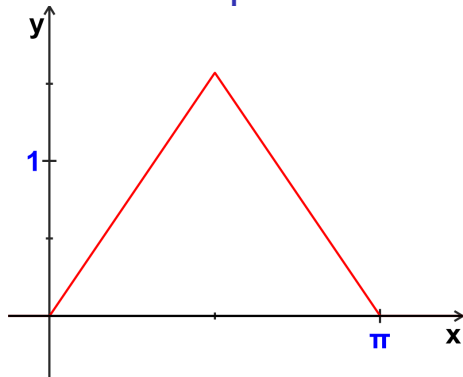
1. $u(x, t)$ is the vertical displacement of the particle at position x on the string, at time t .
2. Partial differential equation: $\frac{\partial^2}{\partial t^2}u(x, t) = \frac{\partial^2}{\partial x^2}u(x, t)$.
(We're assuming $k = 1$.)
3. Initial condition. For $0 \leq x \leq \pi$: $u(x, 0) = f(x)$ (see next slide), $\frac{\partial}{\partial t}u(x, 0) = 0$ (string is not moving as it is released).
4. $u(0, t) = u(\pi, t) = 0$ for $t > 0$ (endpoints of the string are fixed).

The Initial Shape

The Initial Shape



The Initial Shape



(We'll need to think about how to encode it.)

Separating the Equation

Separating the Equation

$$\frac{\partial^2}{\partial t^2}u(x,t) = \frac{\partial^2}{\partial x^2}u(x,t)$$

Separating the Equation

$$\frac{\partial^2}{\partial t^2} u(x, t) = \frac{\partial^2}{\partial x^2} u(x, t) \quad u(x, t) := X(x)T(t)$$

Separating the Equation

$$\begin{aligned}\frac{\partial^2}{\partial t^2}u(x,t) &= \frac{\partial^2}{\partial x^2}u(x,t) & u(x,t) &:= X(x)T(t) \\ X(x)T''(t) &= X''(x)T(t)\end{aligned}$$

Separating the Equation

$$\begin{aligned}\frac{\partial^2}{\partial t^2}u(x,t) &= \frac{\partial^2}{\partial x^2}u(x,t) & u(x,t) &:= X(x)T(t) \\ X(x)T''(t) &= X''(x)T(t) \\ \frac{T''(t)}{T(t)} &= \frac{X''(x)}{X(x)}\end{aligned}$$

Separating the Equation

$$\begin{aligned}\frac{\partial^2}{\partial t^2}u(x,t) &= \frac{\partial^2}{\partial x^2}u(x,t) & u(x,t) &:= X(x)T(t) \\ X(x)T''(t) &= X''(x)T(t) \\ \frac{T''(t)}{T(t)} &= \frac{X''(x)}{X(x)} = \lambda\end{aligned}$$

Separating the Equation

$$\begin{aligned}\frac{\partial^2}{\partial t^2}u(x,t) &= \frac{\partial^2}{\partial x^2}u(x,t) & u(x,t) &:= X(x)T(t) \\ X(x)T''(t) &= X''(x)T(t) \\ \frac{T''(t)}{T(t)} &= \frac{X''(x)}{X(x)} = \lambda\end{aligned}$$

$$\frac{T''(t)}{T(t)} = \lambda$$

Separating the Equation

$$\begin{aligned}\frac{\partial^2}{\partial t^2}u(x,t) &= \frac{\partial^2}{\partial x^2}u(x,t) & u(x,t) &:= X(x)T(t) \\ X(x)T''(t) &= X''(x)T(t) \\ \frac{T''(t)}{T(t)} &= \frac{X''(x)}{X(x)} = \lambda\end{aligned}$$

$$\frac{T''(t)}{T(t)} = \lambda \quad \text{and} \quad \frac{X''(x)}{X(x)} = \lambda$$

Separating the Equation

$$\begin{aligned}\frac{\partial^2}{\partial t^2}u(x,t) &= \frac{\partial^2}{\partial x^2}u(x,t) & u(x,t) &:= X(x)T(t) \\ X(x)T''(t) &= X''(x)T(t) \\ \frac{T''(t)}{T(t)} &= \frac{X''(x)}{X(x)} = \lambda\end{aligned}$$

$$\frac{T''(t)}{T(t)} = \lambda \quad \text{and} \quad \frac{X''(x)}{X(x)} = \lambda$$

$$T''(t) - \lambda T(t) = 0$$

Separating the Equation

$$\begin{aligned}\frac{\partial^2}{\partial t^2}u(x,t) &= \frac{\partial^2}{\partial x^2}u(x,t) & u(x,t) &:= X(x)T(t) \\ X(x)T''(t) &= X''(x)T(t) \\ \frac{T''(t)}{T(t)} &= \frac{X''(x)}{X(x)} = \lambda\end{aligned}$$

$$\begin{aligned}\frac{T''(t)}{T(t)} = \lambda & \quad \text{and} \quad \frac{X''(x)}{X(x)} = \lambda \\ T''(t) - \lambda T(t) = 0 & \quad \text{and} \quad X''(x) - \lambda X(x) = 0\end{aligned}$$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0)$$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0) = k_1 + k_2$$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0) = k_1 + k_2$$

$$0 = X(\pi)$$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0) = k_1 + k_2$$

$$0 = X(\pi) = k_1 e^{\mu\pi} + k_2 e^{-\mu\pi}$$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0) = k_1 + k_2$$

$$0 = X(\pi) = k_1 e^{\mu\pi} + k_2 e^{-\mu\pi}$$

$$0 = k_1 (e^{\mu\pi} - e^{-\mu\pi}),$$

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0) = k_1 + k_2$$

$$0 = X(\pi) = k_1 e^{\mu\pi} + k_2 e^{-\mu\pi}$$

$$0 = k_1 (e^{\mu\pi} - e^{-\mu\pi}),$$

which forces $k_1 = k_2 = 0$ and then $X(x) = 0$, which cannot be.

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0) = k_1 + k_2$$

$$0 = X(\pi) = k_1 e^{\mu\pi} + k_2 e^{-\mu\pi}$$

$$0 = k_1 (e^{\mu\pi} - e^{-\mu\pi}),$$

which forces $k_1 = k_2 = 0$ and then $X(x) = 0$, which cannot be.

So $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$ with $\mu = \sqrt{|\lambda|}$.

$$X''(x) - \lambda X(x) = 0, X(0) = X(\pi) = 0$$

The solution is either $X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ or it is $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$.

$X(x) = k_1 e^{\mu x} + k_2 e^{-\mu x}$ is not possible because of the boundary conditions:

$$0 = X(0) = k_1 + k_2$$

$$0 = X(\pi) = k_1 e^{\mu\pi} + k_2 e^{-\mu\pi}$$

$$0 = k_1 (e^{\mu\pi} - e^{-\mu\pi}),$$

which forces $k_1 = k_2 = 0$ and then $X(x) = 0$, which cannot be. So $X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x)$ with $\mu = \sqrt{|\lambda|}$. In some presentations, $-\lambda$ is used instead of λ , because the outcome of this computation is anticipated.

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$0 = X(0)$$

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$0 = X(0) = k_1 \cdot 1 + k_2 \cdot 0$$

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$0 = X(0) = k_1 \cdot 1 + k_2 \cdot 0 = k_1$$

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$\begin{aligned} 0 &= X(0) = k_1 \cdot 1 + k_2 \cdot 0 = k_1 \\ X(x) &= k_2 \sin(\mu x) \end{aligned}$$

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$0 = X(0) = k_1 \cdot 1 + k_2 \cdot 0 = k_1$$

$$X(x) = k_2 \sin(\mu x)$$

$$0 = X(\pi)$$

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$0 = X(0) = k_1 \cdot 1 + k_2 \cdot 0 = k_1$$

$$X(x) = k_2 \sin(\mu x)$$

$$0 = X(\pi) = k_2 \sin(\mu \pi)$$

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$0 = X(0) = k_1 \cdot 1 + k_2 \cdot 0 = k_1$$

$$X(x) = k_2 \sin(\mu x)$$

$$0 = X(\pi) = k_2 \sin(\mu \pi)$$

Requires μ to be a nonnegative integer n . (Animation.)

$$X(x) = k_1 \cos(\mu x) + k_2 \sin(\mu x), X(0) = X(\pi) = 0$$

$$\begin{aligned} 0 &= X(0) = k_1 \cdot 1 + k_2 \cdot 0 = k_1 \\ X(x) &= k_2 \sin(\mu x) \\ 0 &= X(\pi) = k_2 \sin(\mu \pi) \end{aligned}$$

Requires μ to be a nonnegative integer n . (Animation.)

So $X(x) = k_1 \sin(nx)$ and $\lambda = -\mu^2 = -n^2$.

Back to T

Back to T

$$\frac{T''}{T} = \lambda$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$
$$T'' + n^2T = 0$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$

$$T'' + n^2 T = 0$$

$$T(t) = c_1 \cos(nt) + c_2 \sin(nt)$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$

$$T'' + n^2T = 0$$

$$T(t) = c_1 \cos(nt) + c_2 \sin(nt)$$

$$T'(t) = -c_1 n \sin(nt) + c_2 n \cos(nt)$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$

$$T'' + n^2T = 0$$

$$T(t) = c_1 \cos(nt) + c_2 \sin(nt)$$

$$T'(t) = -c_1 n \sin(nt) + c_2 n \cos(nt)$$

$$0 = T'(0) \quad \text{initial condition}$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$

$$T'' + n^2T = 0$$

$$T(t) = c_1 \cos(nt) + c_2 \sin(nt)$$

$$T'(t) = -c_1 n \sin(nt) + c_2 n \cos(nt)$$

$$0 = T'(0) \quad \text{initial condition}$$

$$0 = -c_1 n \cdot 0 + c_2 n \cdot 1$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$

$$T'' + n^2T = 0$$

$$T(t) = c_1 \cos(nt) + c_2 \sin(nt)$$

$$T'(t) = -c_1 n \sin(nt) + c_2 n \cos(nt)$$

$$0 = T'(0) \quad \text{initial condition}$$

$$0 = -c_1 n \cdot 0 + c_2 n \cdot 1$$

$$0 = c_2$$

Back to T

$$\frac{T''}{T} = \lambda = -n^2$$

$$T'' + n^2T = 0$$

$$T(t) = c_1 \cos(nt) + c_2 \sin(nt)$$

$$T'(t) = -c_1 n \sin(nt) + c_2 n \cos(nt)$$

$$0 = T'(0) \quad \text{initial condition}$$

$$0 = -c_1 n \cdot 0 + c_2 n \cdot 1$$

$$0 = c_2$$

So $T(t) = c_1 \cos(nt)$.

Back to u

Back to u

$$u(x, t) = X(x)T(t)$$

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)\end{aligned}$$

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt)\end{aligned}$$

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt) \\ u_n(x, t) &= b_n \sin(nx) \cos(nt)\end{aligned}$$

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt)\end{aligned}$$

$$u_n(x, t) = b_n \sin(nx) \cos(nt)$$

None of these solutions fit our initial condition.

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt)\end{aligned}$$

$$u_n(x, t) = b_n \sin(nx) \cos(nt)$$

None of these solutions fit our initial condition. But because all u_n solve the equation,

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt)\end{aligned}$$

$$u_n(x, t) = b_n \sin(nx) \cos(nt)$$

None of these solutions fit our initial condition. But because all u_n solve the equation, the boundary condition

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt)\end{aligned}$$

$$u_n(x, t) = b_n \sin(nx) \cos(nt)$$

None of these solutions fit our initial condition. But because all u_n solve the equation, the boundary condition and the initial condition on the time derivative,

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt)\end{aligned}$$

$$u_n(x, t) = b_n \sin(nx) \cos(nt)$$

None of these solutions fit our initial condition. But because all u_n solve the equation, the boundary condition and the initial condition on the time derivative, so should

$$u(x, t) = \sum_{n=1}^{\infty} b_n u_n(x, t) = \sum_{n=1}^{\infty} b_n \sin(nx) \cos(nt).$$

Back to u

$$\begin{aligned}u(x, t) &= X(x)T(t) \\ &= k_1 \sin(nx)c_1 \cos(nt) \\ u_n(x, t) &= b_n \sin(nx) \cos(nt)\end{aligned}$$

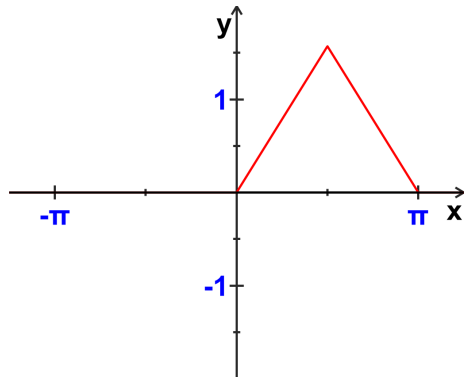
None of these solutions fit our initial condition. But because all u_n solve the equation, the boundary condition and the initial condition on the time derivative, so should

$$u(x, t) = \sum_{n=1}^{\infty} b_n u_n(x, t) = \sum_{n=1}^{\infty} b_n \sin(nx) \cos(nt).$$

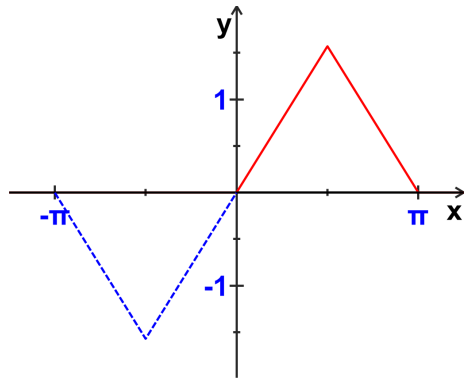
(Significant theory required to assure the infinite summation does not destroy anything.)

A Trick to Allow the Use of Sine Series

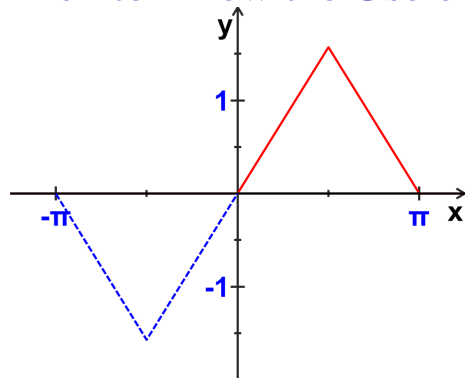
A Trick to Allow the Use of Sine Series



A Trick to Allow the Use of Sine Series



A Trick to Allow the Use of Sine Series



$$f(x) = \begin{cases} \frac{\pi}{2} - \left| x - \frac{\pi}{2} \right|; & \text{for } 0 \leq x \leq \pi, \\ \left| x + \frac{\pi}{2} \right| - \frac{\pi}{2}; & \text{for } -\pi \leq x \leq 0, \end{cases}$$

Fourier Coefficients

Fourier Coefficients

$$b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) dx$$

Fourier Coefficients

$$\begin{aligned} b_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) \, dx \\ &= \frac{2}{\pi} \int_0^{\pi} \left(\frac{\pi}{2} - \left| x - \frac{\pi}{2} \right| \right) \sin(nx) \, dx \end{aligned}$$

Fourier Coefficients

$$\begin{aligned}b_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) \, dx \\ &= \frac{2}{\pi} \int_0^{\pi} \left(\frac{\pi}{2} - \left| x - \frac{\pi}{2} \right| \right) \sin(nx) \, dx \\ b_{2k} &= 0,\end{aligned}$$

because for even n , the integrand is odd with respect to the center $\frac{\pi}{2}$.

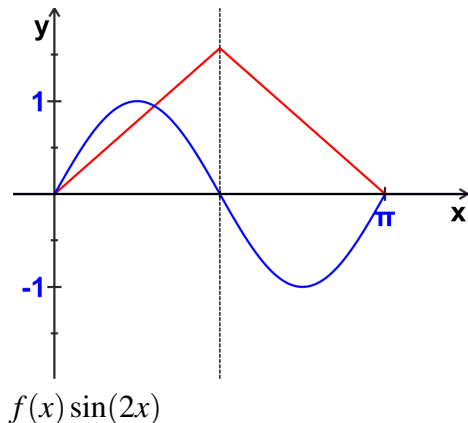
Fourier Coefficients

$$\begin{aligned}b_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) \, dx \\ &= \frac{2}{\pi} \int_0^{\pi} \left(\frac{\pi}{2} - \left| x - \frac{\pi}{2} \right| \right) \sin(nx) \, dx \\ b_{2k} &= 0,\end{aligned}$$

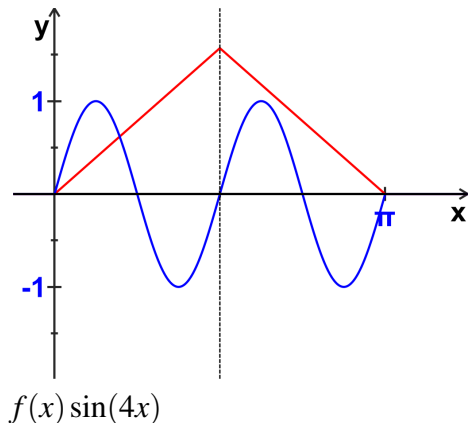
because for even n , the integrand is odd with respect to the center $\frac{\pi}{2}$. For odd n , the integrand is even with respect to the center $\frac{\pi}{2}$ and we continue as follows.

Even/Odd With Respect to $\frac{\pi}{2}$

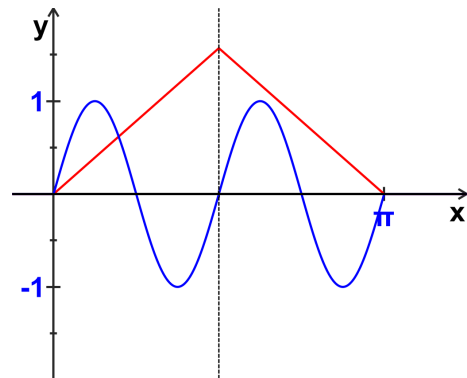
Even/Odd With Respect to $\frac{\pi}{2}$



Even/Odd With Respect to $\frac{\pi}{2}$

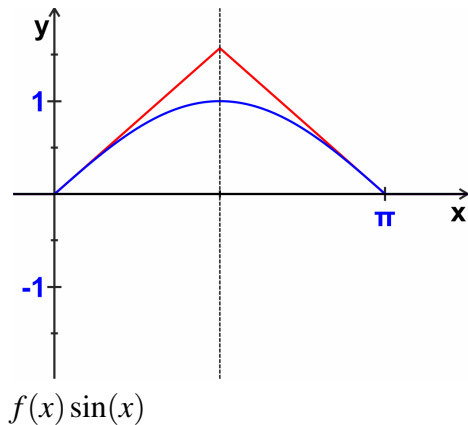


Even/Odd With Respect to $\frac{\pi}{2}$

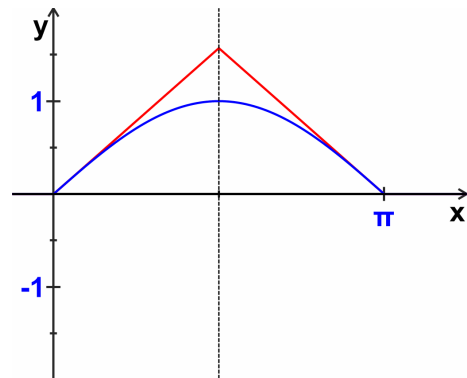


$f(x) \sin(4x)$
(and so on)

Even/Odd With Respect to $\frac{\pi}{2}$

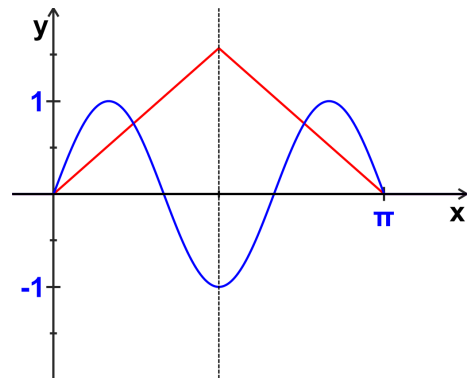


Even/Odd With Respect to $\frac{\pi}{2}$



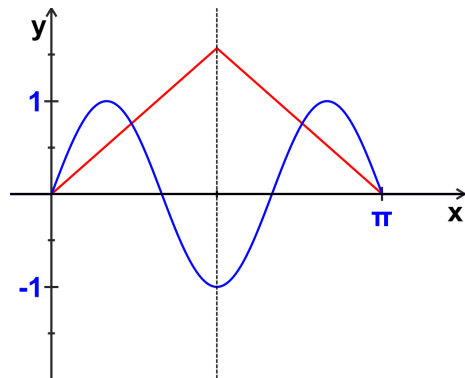
$f(x)\sin(x)$
(Doubles)

Even/Odd With Respect to $\frac{\pi}{2}$



$f(x) \sin(3x)$
(Doubles)

Even/Odd With Respect to $\frac{\pi}{2}$



$f(x) \sin(3x)$
 (Doubles)
 (and so on)

Fourier Coefficients

Fourier Coefficients

$$b_{2k+1} = \frac{4}{\pi} \int_0^{\frac{\pi}{2}} \left(\frac{\pi}{2} - \left| x - \frac{\pi}{2} \right| \right) \sin((2k+1)x) dx$$

Fourier Coefficients

$$\begin{aligned} b_{2k+1} &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} \left(\frac{\pi}{2} - \left| x - \frac{\pi}{2} \right| \right) \sin((2k+1)x) \, dx \\ &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} \left[\frac{\pi}{2} - \left(\frac{\pi}{2} - x \right) \right] \sin((2k+1)x) \, dx \end{aligned}$$

Fourier Coefficients

$$\begin{aligned}b_{2k+1} &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} \left(\frac{\pi}{2} - \left| x - \frac{\pi}{2} \right| \right) \sin((2k+1)x) \, dx \\ &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} \left[\frac{\pi}{2} - \left(\frac{\pi}{2} - x \right) \right] \sin((2k+1)x) \, dx \\ &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} x \sin((2k+1)x) \, dx\end{aligned}$$

Fourier Coefficients

$$b_{2k+1} = \frac{4}{\pi} \int_0^{\frac{\pi}{2}} x \sin((2k+1)x) dx$$

Fourier Coefficients

$$\begin{aligned} b_{2k+1} &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} x \sin((2k+1)x) dx \\ &= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) \right]_0^{\frac{\pi}{2}} \end{aligned}$$

Fourier Coefficients

$$\begin{aligned} b_{2k+1} &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} x \sin((2k+1)x) dx \\ &= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) \Big|_0^{\frac{\pi}{2}} + \frac{1}{2k+1} \int_0^{\frac{\pi}{2}} \cos((2k+1)x) dx \right] \end{aligned}$$

Fourier Coefficients

$$\begin{aligned} b_{2k+1} &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} x \sin((2k+1)x) dx \\ &= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) \Big|_0^{\frac{\pi}{2}} + \frac{1}{2k+1} \int_0^{\frac{\pi}{2}} \cos((2k+1)x) dx \right] \\ &= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) + \frac{1}{(2k+1)^2} \sin((2k+1)x) \right]_0^{\frac{\pi}{2}} \end{aligned}$$

Fourier Coefficients

$$\begin{aligned}b_{2k+1} &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} x \sin((2k+1)x) dx \\&= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) \Big|_0^{\frac{\pi}{2}} + \frac{1}{2k+1} \int_0^{\frac{\pi}{2}} \cos((2k+1)x) dx \right] \\&= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) + \frac{1}{(2k+1)^2} \sin((2k+1)x) \right]_0^{\frac{\pi}{2}} \\&= -\frac{2}{2k+1} \cos\left((2k+1)\frac{\pi}{2}\right) + \frac{4}{\pi} \frac{1}{(2k+1)^2} \sin\left((2k+1)\frac{\pi}{2}\right)\end{aligned}$$

Fourier Coefficients

$$\begin{aligned}b_{2k+1} &= \frac{4}{\pi} \int_0^{\frac{\pi}{2}} x \sin((2k+1)x) dx \\&= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) \Big|_0^{\frac{\pi}{2}} + \frac{1}{2k+1} \int_0^{\frac{\pi}{2}} \cos((2k+1)x) dx \right] \\&= \frac{4}{\pi} \left[-x \frac{1}{2k+1} \cos((2k+1)x) + \frac{1}{(2k+1)^2} \sin((2k+1)x) \right]_0^{\frac{\pi}{2}} \\&= -\frac{2}{2k+1} \cos\left((2k+1)\frac{\pi}{2}\right) + \frac{4}{\pi} \frac{1}{(2k+1)^2} \sin\left((2k+1)\frac{\pi}{2}\right) \\&= \frac{4}{\pi} \frac{(-1)^k}{(2k+1)^2}\end{aligned}$$

The Solution

The Solution

$$u(x, t) := \sum_{k=0}^{\infty} \frac{4}{\pi} \frac{(-1)^k}{(2k+1)^2} \sin((2k+1)x) \cos((2k+1)t)$$

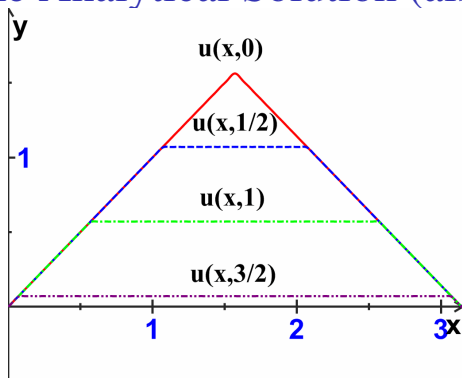
The Solution

$$u(x, t) := \sum_{k=0}^{\infty} \frac{4}{\pi} \frac{(-1)^k}{(2k+1)^2} \sin((2k+1)x) \cos((2k+1)t)$$

But what does that look like?

The Analytical Solution (also animated)

The Analytical Solution (also animated)



The Real Experiment

The Real Experiment

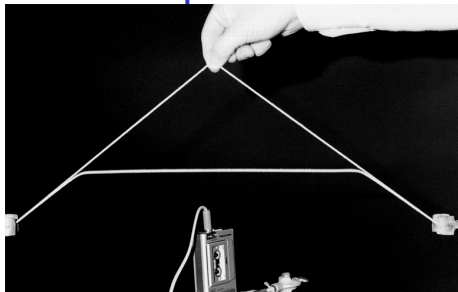


Image courtesy of Loren M. Winters, used with permission.

The Real Experiment

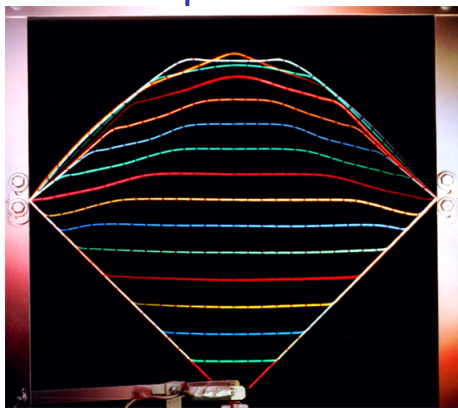


Image courtesy of Loren M. Winters, used with permission.